3) Suppose that the group $G$ has a representation $G=\left\langle a, b, c, d \mid R_{i}=e, i=1-n\right\rangle$ with a finite set of defining relations. We call the quantity $\Sigma \partial\left(R_{i}\right)$ the length of the corresponding representation. Among all such representations, we choose one with minimal length. The representation of $H$ (as a subgroup), constructed using the minimal representation of $G$, has lesser length than the original representation of $G$. Therefore, the factor-representation of $G$, obtained from the representation of $H$ by applying the homomorphism described in 1), has lesser length than the original minimal representation of $G$. This contradiction shows that the group $G$ cannot be defined by a finite set of defining relations. The theorem is proved.

Despite the fact that the group $G$ was defined as a group of transformations on a space with a measure, it may be defined in purely algebraic terms. In particular, there exists a simple algorithm, allowing us to answer the following question for any word $W$ in the generators of the group $G$ : does $W$ represent the unit element of $G$, or not? We also note that $G$ is finitely approximable and has exponential growth. We give another example. Let $\xi$ be the transformation of the square $[0,1] \times[0,1]$, consisting of the cyclic permutation of its quadrants, and let the transformation $\eta$ be described as in Fig. 2 ( $S$ denotes the cyclic permutation of the quadrants of the square over which it is written). Then $\xi^{4}=\eta^{4}=e$ and the group generated by the transformations $\xi$ and $\eta$ is an infinite periodic group.

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## VARIATIONAL PRINCIPLE FOR EQUATIONS INTEGRABLE

BY THE INVERSE PROBLEM METHOD
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We consider (cf. [1]) the system of nonlinear equations representing the conditions for compatibility of two linear differential equations for a square $\mathrm{N} \times \mathrm{N}$ nonsingular matrix function $\Psi(\xi, \eta, \lambda)$

$$
\begin{equation*}
\Psi_{\xi}=U(\xi, \eta, \lambda) \Psi, \quad \Psi_{\eta}=V(\xi, \eta, \lambda) \Psi \tag{1}
\end{equation*}
$$

Here $U$ and $V$ are rational functions of the parameter $\lambda$ with distinct simple poles:

$$
\begin{equation*}
U=U_{0}+\sum_{n=1}^{N_{1}} \frac{U_{n}(\xi, \eta)}{\lambda-a_{n}}, \quad V=V_{0}+\sum_{n=1}^{N_{2}} \frac{V_{n}(\xi, \eta)}{\lambda-b_{n}} . \tag{2}
\end{equation*}
$$

The compatibility conditions for Eqs. (1) have the form

$$
\begin{gather*}
U_{0 \eta}-V_{0 \xi}=\left[U_{0}, V_{0}\right],  \tag{3}\\
U_{n \eta}=\left[U_{n}, V_{0}+\sum_{k=1}^{N_{2}} \frac{V_{k}}{a_{n}-b_{k}}\right], \quad V_{n \mathrm{~s}}=\left[V_{n}, U_{0}+\sum_{k=1}^{N_{1}} \frac{U_{k}}{b_{n}-a_{k}}\right] . \tag{4}
\end{gather*}
$$

It follows from (3) that there exists a nonsingular matrix $g(\xi, \eta)$ such that

$$
\begin{equation*}
U_{0}=g \xi g^{-1}, \quad V_{0}=g_{\eta} g^{-1} . \tag{5}
\end{equation*}
$$

We introduce the notation

$$
\nabla_{\eta}=\frac{\partial}{\partial \eta}-V_{0}, \quad \nabla_{\xi}=\frac{\partial}{\partial \xi}-U_{0} .
$$

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We express the matrices $U_{n}, V_{n}$ in the form

$$
\begin{equation*}
U_{n}=\varphi_{n} V_{n}^{(0)} \varphi_{n}^{-1}, \quad V_{n}=\psi_{n} V_{n}^{(0)} \psi_{n}^{-1}, \tag{6}
\end{equation*}
$$

where the $U_{n}^{(0)}(\xi), V_{n}^{(0)}(\eta)$ are the normal Jordan forms of $U_{n}, V_{n}$. They are obtained by partial integration of Eqs. (4). Substituting (6) into (3), (4), we find the equations satisfied by the matrices $\varphi_{\mathrm{n}}, \psi_{\mathrm{n}}$,

$$
\begin{gather*}
\left(\nabla_{\eta} \varphi_{n}-\sum_{m=1}^{N_{2}} \frac{\psi_{m} V_{m}^{(0)} \psi_{m}^{-1}}{a_{n}-b_{m}} \varphi_{n}\right) U_{n}^{(0)}=U_{n}\left(\nabla_{\eta} \varphi_{n}-\sum_{m=1}^{N_{\ddagger}} \frac{\psi_{m} V_{m}^{(0)} \psi_{m}^{-1}}{a_{n}-b_{m}} \varphi_{n}\right),  \tag{7}\\
\left(\nabla_{\xi} \psi_{n}-\sum_{m=1}^{N_{1}} \frac{\varphi_{m} U_{m}^{(0)} \varphi_{m}}{b_{n}-a_{m}} \psi_{n}\right) V_{n}^{(0)}=v_{n}\left(\nabla_{\xi} \psi_{n}-\sum_{m=1}^{N_{1}} \frac{\varphi_{m} U_{m}^{(0)} \varphi_{m}^{-1}}{b_{n}-a_{m}} \psi_{n}\right) .
\end{gather*}
$$

Equations (7) can be rewritten as

$$
\begin{equation*}
\nabla_{\eta} \varphi_{n}-\sum_{m=1}^{N_{1}} \frac{\psi_{m} V_{m}^{(0)} \psi_{m}^{-1}}{a_{n}-b_{m}} \varphi_{n}=\varphi_{n} 4_{n}, \quad \nabla_{\xi} \psi_{n}-\sum_{m=1}^{N_{n}} \frac{\varphi_{m} U_{m}^{(0)} \varphi_{m}^{-1}}{b_{n}-a_{m}} \psi_{n}=\psi_{n} B_{n} \tag{8}
\end{equation*}
$$

Here $A_{n}(\xi, \eta), B_{n}(\xi, \eta)$ are arbitrary matrix functions commuting with the matrices $U_{n}^{(0)}$ and $V_{n}^{(0)}$, respectively. Their appearance in Eqs. (8) is related to the obvious nonuniqueness in determining the matrices $\varphi_{\mathrm{n}}, \psi_{\mathrm{n}}$.

Equations (8) imply the easily verified relation

$$
\begin{equation*}
\left[\nabla_{n}, \sum_{n=1}^{N_{1}} U_{n}\right]=\left[\nabla_{5}, \sum_{n=1}^{N_{n}} V_{n}\right] \tag{9}
\end{equation*}
$$

We consider the functional

$$
\begin{equation*}
S=\int_{\xi=1}^{E} d \zeta \int_{\eta_{1}}^{\eta_{2}} d \eta S p\left[\sum_{n=1}^{N_{1}} \varphi_{n}^{-1} \nabla_{\eta} \varphi_{n} U_{n}^{(0)}-\sum_{n=1}^{N_{1}} \psi_{n}^{-1} \nabla_{\xi} \psi_{n}{ }_{n}^{(1)}-\sum_{n=1}^{N_{1}} \sum_{m=1}^{N_{2}} \frac{\psi_{m_{m}} V_{m}^{(0)} \psi_{m}^{-1} \varphi_{n} U_{n}^{(0)} \varphi_{n}^{-1}}{a_{n}-b_{m}}\right] \tag{10}
\end{equation*}
$$

where $\xi_{1}, \xi_{2}, \eta_{1}, \eta_{2}$ are chosen arbitrarily, and find conditions such that the variation $\delta S$ vanishes. It is easy to check that variation with respect to $\varphi_{\mathrm{n}}, \psi_{\mathrm{n}}$ leads to Eqs. (7), while variation with respect to $g$ leads to Eqs. (9).

Thus, the functional $S$ is an action for Eqs. (7). If $S$ is real, then Eqs. (7) are a Hamiltonian system for which the Hamiltonian and symplectic form are evaluated in the obvious way. In general, the question of whether the Hamiltonian structure constructed is unique remains open.

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