

Resonant control of solitons

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Introduction

Autoresonant excitation of solitons

KdV: I.Aranson, B.Meerson, T.Tajima, *Phys.Rev.A* 45 (1992) 7500

NLS: L.Friedland, A.Shagalov, *Phys.Rev.Lett.* 81 (1998) 4357

Autoresonant control (in particular, amplification) of solitons

KdV: R.Grimshow, E.Pelinovsky, P.Sakov,
Studies in Applied Mathematics 97 (1996) 235

SG: E.Maslov, L.Kalykin, A.Shagalov, *Theor.Math.Phys.* 152 (2007) 356

NLS: S.Batalov, E.Maslov, A.Shagalov, *Sov. Phys. JEPT* 108 (2009) 890

Problems of the autoresonant control of solitons:

- (1) narrow range of the driving frequencies
- (2) strong restrictions on the phase of the driving

Resonant perturbation

$$iu_t + \frac{1}{2}u_{xx} + |u|^2u = \varepsilon e^{i\psi(t)}, \quad 0 \leq \varepsilon \ll 1, \quad (NLS)$$

with a slowly varying frequency: $\Omega(t) = \psi_t(t), \quad \Omega_t \sim \varepsilon$

Soliton solution ($\varepsilon = 0$): $u_s = \frac{ae^{i\theta}}{\cosh[a(x - \xi_0)]}, \quad \theta = \omega t + \theta_0,$

soliton amplitude: a

soliton frequency: $\omega = a^2/2$

Resonance of the soliton and the driving:

$$\omega \approx \Omega(t)$$

Equations for soliton parameters

$$\delta_t = \Delta\Omega - \frac{\varepsilon\pi a}{\Omega} \cos\delta$$

$$a_t = -\frac{\varepsilon\pi a^2}{2\Omega} \sin\delta$$

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JETP 108, 1021 (2009)*

$$\delta = \theta - \psi \quad \text{— phase shift between the soliton and the driving}$$

$$\Delta\Omega(t) = a^2/2 - \Omega(t) \quad \text{— difference in frequencies of the soliton and the driving}$$

suppose: $\Omega(t) = \Omega(0) - \alpha t$, $\alpha \sim \varepsilon$, $\Delta\Omega \sim \sqrt{\varepsilon}$

introduce "slow" time: $\tau = \sqrt{\varepsilon} t$.

Equations for soliton parameters

$$\delta(\tau) = \delta_0(\tau) + \sqrt{\varepsilon}\delta_1(\tau) + \varepsilon\delta_2(\tau) + \dots$$

$$a(\tau) = a_0 + \sqrt{\varepsilon}a_1(\tau) + \varepsilon a_2(\tau) + \dots$$

$$a(0) = a_0$$

$$\Omega = \underbrace{\Omega_0 + \sqrt{\varepsilon}\Omega_1}_{\Omega(0)} - \sqrt{\varepsilon}\beta\tau, \quad \Omega_0 = a_0^2/2, \quad \beta = \alpha/\varepsilon \sim 1$$

one finds in the lowest order

$$\delta_{0,\tau} = a_0 a_1 - \Omega_1 + \beta\tau,$$

$$a_{1,\tau} = -\pi \sin \delta_0$$

equation for the phase difference:

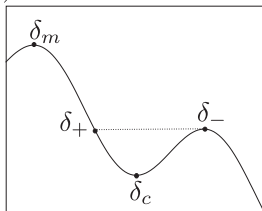
$$\delta_{0,\tau\tau} = -\pi a_0 \sin \delta_0 + \beta$$

Scattering on resonance

$$V(\delta_0) = -\pi a_0 \cos \delta_0 - \beta \delta_0,$$

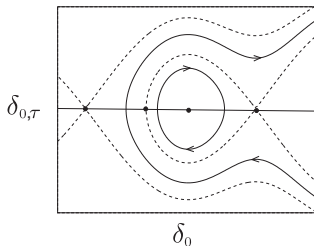
$$\lambda = \frac{\beta}{\pi a_0}$$

$V(\delta_0)$



$0 < \lambda < 1 \implies$ threshold:

$$\varepsilon > \varepsilon_{th} = \frac{\alpha}{\pi a_0}$$

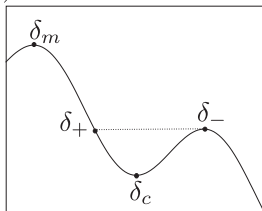


Scattering on resonance

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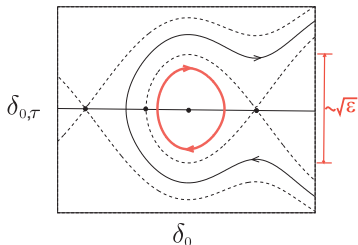
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$V(\delta_0)$



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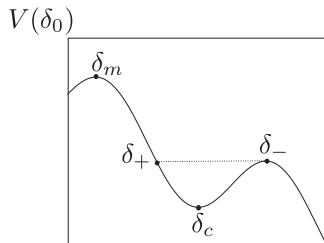
$$\varepsilon > \varepsilon_{th} = \frac{\alpha}{\pi a_0}$$



$$\Delta\Omega(o) \sim \sqrt{\varepsilon}, \quad \delta_+ < \delta(o) < \delta_-$$

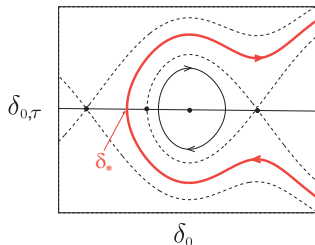
Scattering on resonance

$$V(\delta_0) = -\pi a_0 \cos \delta_0 - \beta \delta_0, \quad \lambda = \frac{\beta}{\pi a_0}$$



$0 < \lambda < 1 \implies$ threshold:

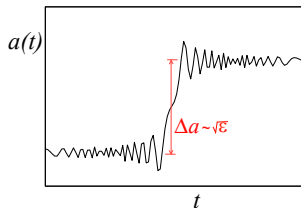
$$\varepsilon > \varepsilon_{th} = \frac{\alpha}{\pi a_0}$$



Resonance: $\delta_{0,\tau} \approx 0$

Scattering on resonance

Normal point: $\delta_0(0) = \delta_*$ ($\delta_m < \delta_* < \delta_+$)



$$\Delta a = -\pi\sqrt{\varepsilon} \int_{-\infty}^{\infty} \sin \delta_0(\tau) d\tau$$

$$\delta_0(\tau) \approx \delta^* - C\tau^2/2,$$

$$C = \pi a_0 \sin \delta^* - \beta$$

$$(C \rightarrow 0, \quad \delta_* \rightarrow \delta_m)$$

$$(\Delta a)_0 = -\frac{\sqrt{\varepsilon}\pi^{3/2}(\text{sign}(\alpha) \cos \delta^* + \sin \delta^*)}{\sqrt{|C|}}$$

similar approach: B.V.Chirikiv (1959),

A.I.Neishtadt (1975)

Scattering on resonance

Singular point: $\delta_* \approx \delta_+$

near the saddle point $\delta_0 = \delta_-$ ($\delta_{0,\tau} > 0$) suppose

$$\delta_0 = \delta_- + \xi(\tau), \quad \xi(0) = 0.$$

$$\xi(\tau) = \frac{\delta_{0,\tau}(0)}{2\kappa} (e^{\kappa\tau} - e^{-\kappa\tau}), \quad \kappa^2 = \pi a_0 \sqrt{1 - \lambda^2}.$$

$$(\Delta a)_1 = \sqrt{\varepsilon} \frac{2\pi\lambda}{\kappa} \mathbf{K}_0 \left(\frac{\delta_{0,\tau}(0)}{\kappa} \right),$$

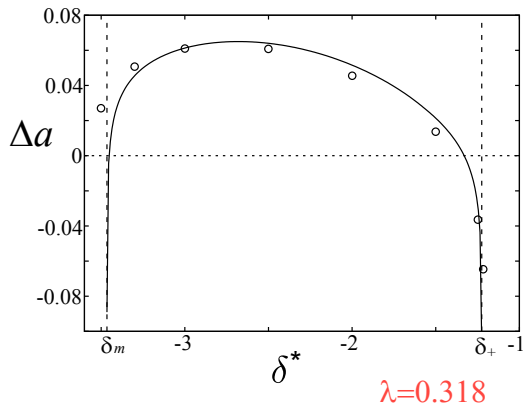
Singular point: $\delta_* \approx \delta_m$

$$\begin{aligned} (\Delta a)_2 = \sqrt{\varepsilon} \frac{\pi^2}{\sqrt{A}} & \left[\sin \left(\delta^* + \frac{C}{A} \right) \mathbf{N}_0 \left(\frac{|C|}{A} \right) \right. \\ & \left. - \text{sign}(\alpha) \cos \left(\delta^* + \frac{C}{A} \right) \mathbf{J}_0 \left(\frac{|C|}{A} \right) \right], \\ & A = -\pi a_0 \cos \delta^*. \end{aligned}$$

Scattering on resonance

$$\delta_* = -\pi/2 : \quad (\Delta a)_2 = (\Delta a)_0$$

$$\Delta a = \begin{cases} 2(\Delta a)_1 + (\Delta a)_2, & \delta_m < \delta^* < -\pi/2, \\ 2(\Delta a)_1 + (\Delta a)_0, & -\pi/2 < \delta^* < \delta_+, \end{cases}$$



Numerical simulation of soliton amplification in NLS

Multiple scattering for an impulse-like ("saw-tooth") modulation of the driving frequency:

$$\Omega(t) = \Omega_{\max} - \alpha[t - (n - 1)T],$$
$$(n - 1)T \leq t \leq nT, \quad n = 1, 2, 3, \dots$$

n - the number of the impulse

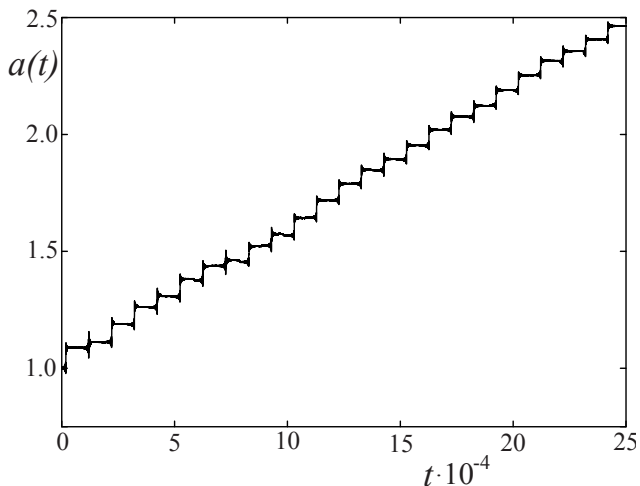
T - its period.

In a single impulse the frequency is changed in the range from Ω_{\max} to $\Omega_{\min} = \Omega_{\max} - \alpha T$

Numerical simulation of soliton amplification in NLS

$$a(0) = 1; \alpha = 0.00025, \varepsilon = 0.0005, \Omega_{\min} = 0.4, \Omega_{\max} = 3, \\ T = 1.04 \cdot 10^4.$$

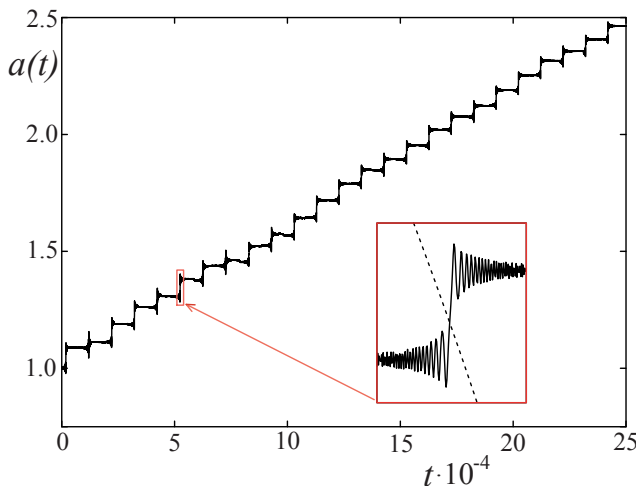
$$\lambda = 0.16$$



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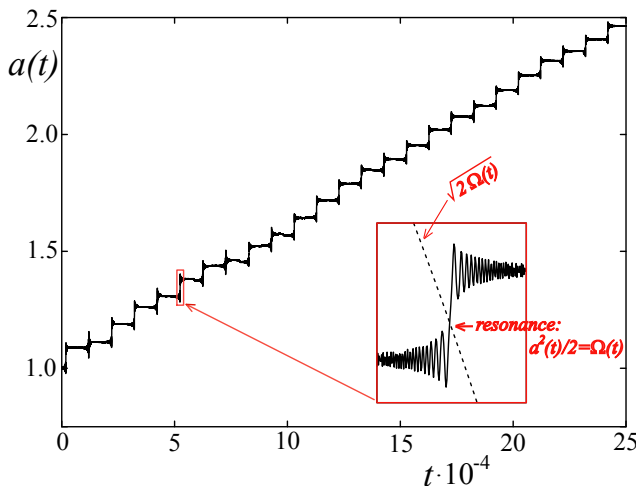
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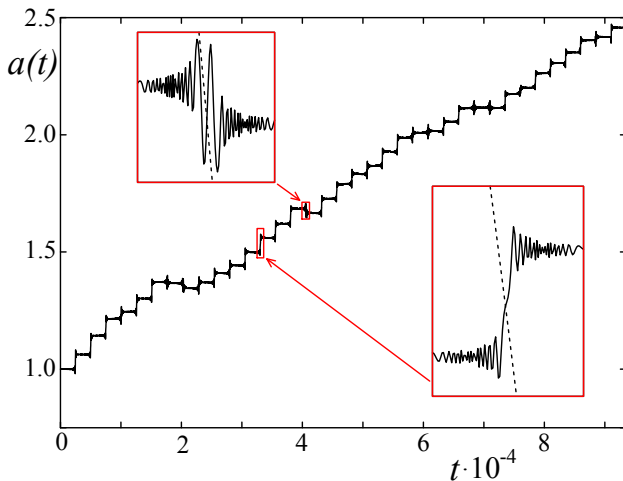
$$\lambda = 0.16$$



Numerical simulation of soliton amplification in NLS

$$a(0) = 1; \alpha = 0.001, \varepsilon = 0.0005, \Omega_{\min} = 0.4, \Omega_{\max} = 3, \\ T = 2.6 \cdot 10^3.$$

$$\lambda = 0.64$$



Control of the soliton in NLS

$$a(0) = 1; \varepsilon = 0.0005, \Omega_{\min} = 0.4, \Omega_{\max} = 3, T = 1.04 \cdot 10^4.$$

